1. (linear partial differential operators and Laplace's equation on a rectangle) Solve the following boundary value problem for  $u \in C^2([0,1] \times [0,2])$ :

$$
\begin{cases}\n\Delta u = 0, \\
u(x, 0) = 0, \ u(1, y) = \sinh \pi \sin \pi y, \ u(x, 2) = \sinh 6\pi \sin 3\pi x, \ u(0, y) = 0.\n\end{cases}
$$
\n(1)

Hint: Consider two separate boundary value problems with homogeneous boundary conditions on three of the four boundary segments of the rectangle. Solve each of these problems using separation of variables. Then use the linearity of the Laplace operator.

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2. (first order linear PDE; method of characteristics) Solve the PDE

$$
xu_x - yu_y + (x^2 + y^2)u = x^2 - y^2 \quad \text{on } U = \{(x, y) \in \mathbb{R}^2 : x, y > 0\}.
$$

"Solve" here means "Find all possible  $C<sup>1</sup>$  solutions." Your solution should depend on an arbitrary function which you will need to introduce. Knowing how to do that is part of the problem. (This is like if someone says: Solve  $x'' = 0$ . Then you know  $x = at + b$ with two arbitrary constants  $a$  and  $b$ .)

Hint(s): Consider the **characteristic field v** =  $(x, -y)$  on the first quadrant U. Plot it with numerical software if necessary. Choose an appropriate non-characteristic curve.

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3. (one dimensional wave equation) Solve the initial value problem for the wave equation:

$$
\begin{cases}\n u_{tt} = u_{xx} \text{ on } \mathbb{R} \times [0, \infty) \\
 u(x, 0) = u_0(x) \\
 u_t(x, 0) = v_0(x)\n\end{cases}
$$
\n(2)

where  $u_0 \in C^2(\mathbb{R})$  and  $v_0 \in C^1(\mathbb{R})$  to obtain d'Alembert's solution:

$$
u(x,t) = \frac{1}{2}[u_0(x+t) + u_0(x-t)] + \frac{1}{2}\int_{x-t}^{x+t} v_0(\xi) d\xi.
$$

Hint(s): Factor the operator  $\Box u = u_{tt} - u_{xx}$  as either

$$
(u_t - u_x)_t + (u_t - u_x)_x
$$
 or  $(u_t + u_x)_t - (u_t + u_x)_x$ .

Then solve two first order PDEs with appropriate Cauchy conditions. Incidentally, the initial conditions in (2) are Cauchy conditions for the wave equation.

Name and section:

## The Divergence of a Vector Field in  $\mathbb{R}^2$

This will give you a chance to integrate on curves. You'll need to integrate on curves. So the first part is a warm up involving integration on a curve. Remember, before you start it, that

$$
\int_{\Gamma} f = \lim \sum f(p_j) \operatorname{length}(\Gamma_j)
$$

where  $\{\Gamma_i\}$  is a partition of the curve  $\Gamma$  and  $f : \Gamma \to \mathbb{R}$  is a real valued function; each point  $p_j$  is in the partition piece  $\Gamma_j$  and the limit is as the "diameter measure" (in this case length will work) of the largest partition piece tends to zero.

Also, the divergence for a vector field  $\mathbf{v} : \mathbb{R}^2 \to \mathbb{R}^2$  at a point  $\mathbf{p} \in \mathbb{R}^2$  is defined as

$$
\operatorname{div} \mathbf{v} = \lim_{U \to \{\mathbf{p}\}} \frac{1}{\operatorname{area}(U)} \oint_{\partial U} \mathbf{v} \cdot \mathbf{n}
$$

(when this limit exists). Here **n** is the **outward unit normal** to  $U$  and the little circle is put on the integral sign just to remind us that we're integrating over a cycle or, in this case, a simple closed curve.

Okay, let's do this.

4. A nice curve to consider (when thinking about integrating on a curve) is a single turn of a helix

$$
\Gamma = \{(\cos t, \sin t, t) : t \in [0, 2\pi]\}.
$$

 $\int f$ Γ

Let's try to compute

where  $f = f(x, y, z)$  is just some function I write down. This should illustrate how integration over a curve works in general. The first step is to write down a parameterization for the curve.

- (a) Write down a parameterization  $\gamma : [0, 2\pi] \to \mathbb{R}^3$  for the specific curve  $\Gamma$  given above and sketch the image. (Hint: Yes, this is as easy as it looks.)
- (b) Now, this is perhaps a little harder: For the computation, we want to "change variables" from  $\Gamma$  to  $[0, 2\pi]$ . This requires a scaling factor:

$$
\int_{\Gamma} f = \int_0^{2\pi} f \circ \gamma(t) \, \sigma \, dt.
$$

What is the scaling factor  $\sigma$  for the specific helix parameterized in the previous part? And what is the scaling factor, in general, if a curve  $\Gamma$  is parameterized by  $\gamma \in C^1([a, b] \to \mathbb{R}^n)$  on some interval  $[a, b]$ ?

(c) Compute

$$
\int_{\Gamma} f
$$

for  $\Gamma$  the single turn of the helix above and  $f(x, y, z) = x^2 + y^2 + z^2$ . Hint: This may not be as easy as it looks at first.

(d) Consider a point  $\mathbf{p} = (p_1, p_2) \in \mathbb{R}^2$  and a vector field  $\mathbf{v} \in C^1(\mathbb{R}^2 \to \mathbb{R}^2)$ . For positive numbers  $\epsilon$  and  $\delta$ , let

$$
R = R_{\epsilon,\delta} = \{ \mathbf{x} = (x_1, x_2) \in \mathbb{R}^2 : |x_1 - p_1| < \epsilon \text{ and } |x_2 - p_2| < \delta \}
$$

be a rectangular domain with outward unit normal  $n$ . Draw R along with  $n$  and show

$$
\int_{\partial R} \mathbf{v} \cdot \mathbf{n} = 2\delta \int_{p_1-\epsilon}^{p_1+\epsilon} \frac{\partial v_2}{\partial y}(x, p_2^*) dx + 2\epsilon \int_{p_2-\delta}^{p_2+\delta} \frac{\partial v_1}{\partial x}(p_1^*, y) dy
$$

for some point  $p^* = (p_1^*, p_2^*) \in R$ . Hint: Use the **mean value theorem** which tells you, for example, that if  $v \in C^1(\mathbb{R}^2)$ , then for  $a < b$  and  $y \in \mathbb{R}$ , there is some  $x^* \in (a, b)$  such that

$$
v(b, y) - v(a, y) = \frac{\partial v}{\partial x}(x^*, y) (b - a).
$$

(e) Compute

$$
\lim_{\epsilon,\delta \to 0} \frac{1}{\operatorname{area}(R)} \int_{\partial R} \mathbf{v} \cdot \mathbf{n}.
$$

Name and section:

## The 2D Heat Equation on  $\mathcal{U} \subset \mathbb{R}^2$

5. Derive the heat equation (carefully and from scratch) as it applies to a laminar domain  $U \subset \mathbb{R}^2$ . Start by listing/identifying all the quantities you will use with their units. Let's try this: I'll start you out and give you a sort of outline to follow. When I put an ellipsis  $(\cdots)$ , this will mean there are details for you to fill in—probably lots of them.

quantity identification units θ<sup>2</sup> = θ2(x, y, t), areal or laminar heat energy density [θ2] = [energy] L2 . . . Incidentally, energy has units of work [energy] = [force]L = ML<sup>2</sup> T2 . . . φ~ <sup>2</sup> = φ~ <sup>2</sup>, laminar heat flux field [φ~ <sup>2</sup>] = . . . . . . . . . u = u(x, y, t), temperature [u] = [temperature] Du = Du(x, y, t), temprature gradient [Du] = . . . σ = σ(x, y, u), specific heat capacity [σ] = . . . K<sup>2</sup> = K(x, y, u), laminar thermal conductivity [K2] = . . . . . . . . .

Accounting of rate of change of total energy

$$
\frac{d}{dt} \int_{U} \theta_2 = - \int_{\partial U} \vec{\phi}_2 \cdot \mathbf{n} + \int_{U} Q_2
$$

Law of specific heat ...

Fourier's law . . .

$$
\frac{\partial}{\partial t}[\sigma \rho_2 u] = \text{div}[K_2 Du] + Q_2.
$$

. . .

. . .

Finally, taking  $\sigma \rho_2 = K_2$  (constant) and setting  $f = Q_2/K_2$ ,

$$
u_t = \Delta u + f.
$$