

Assignment 7:  
Wave equation and Sturm Liouville Theory  
Due Tuesday December 7, 2021

John McCuan

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**Problem 1** (*d'Alembert's solution, Haberman 4.4.6-8*) Consider the initial value problem for the 1-D wave equation

$$\begin{cases} u_{tt} = \sigma^2 u_{xx} & \text{on } \mathbb{R} \times (0, \infty) \\ u(x, 0) = u_0(x), & x \in \mathbb{R} \\ u_t(x, 0) = v_0(x), & x \in \mathbb{R} \end{cases}$$

where  $u_0, v_0 \in C^1(\mathbb{R})$  are two given functions.

(a) Let  $Lu = u_t + \sigma u_x$  and  $Mu = u_t - \sigma u_x$ . Use the method of characteristics (twice) applied to the factorization  $\square u = MLu = 0$  to obtain d'Alembert's solution

$$u(x, t) = \frac{1}{2}[u_0(x - \sigma t) + u_0(x + \sigma t)] + \frac{1}{2\sigma} \int_{x - \sigma t}^{x + \sigma t} v_0(\xi) d\xi.$$

(b) Show that if  $u_0$  and  $v_0$  are both odd and periodic with period  $2L$ , then the restriction of d'Alembert's solution to  $[0, L] \times (0, \infty)$  satisfies the initial/boundary value problem

$$\begin{cases} u_{tt} = \sigma^2 u_{xx} & \text{on } \mathbb{R} \times (0, \infty) \\ u(0, t) = 0 = u(L, t), & t \geq 0 \\ u(x, 0) = u_0(x), & x \in \mathbb{R} \\ u_t(x, 0) = v_0(x), & x \in \mathbb{R} \end{cases}$$

(c) Consider the initial/boundary value problem

$$\begin{cases} u_{tt} = \sigma^2 u_{xx} & \text{on } \mathbb{R} \times (0, \infty) \\ u(0, t) = 0 = u(L, t), & t > 0 \\ u(x, 0) = g(x), & x \in \mathbb{R} \\ u_t(x, 0) = h(x), & x \in \mathbb{R} \end{cases}$$

where  $g, h \in C^1[0, L]$ . Can d'Alembert's solution be applied to solve this problem? Why or why not?

**Problem 2** (The heat equation on  $\mathbb{R}$ ; fundamental solution, Haberman section 10.4)  
We have considered special solutions of the heat equation having the form

$$u(x, t) = e^{-j^2 \pi^2 t / L^2} \cos\left(\frac{j\pi}{L} x\right) \quad \text{and} \quad u(x, t) = e^{-j^2 \pi^2 t / L^2} \sin\left(\frac{j\pi}{L} x\right)$$

on the interval  $[0, L]$ . These are **separated variables** solutions. They can, of course, also be considered as solutions on all of the spatial domain  $\mathbb{R}$ , but that consideration is not so interesting because they are spatially periodic. There is another important solution of the heat equation to know about and remember.

The function  $\Phi : \mathbb{R} \times (0, \infty) \rightarrow \mathbb{R}$  given by

$$\Phi(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}}$$

is called the **fundamental solution** of the one-dimensional heat equation. Notice that the fundamental solution does not have the form of a separated variables solution.

(a) Verify that  $\Phi$  satisfies  $u_t = u_{xx}$  for  $(x, t) \in \mathbb{R} \times (0, \infty)$ .

(b) Use L'Hopital's rule to determine

$$\lim_{t \searrow 0} \Phi(x, t).$$

(c) Make an animation of the spatial graph of the fundamental solution  $\Phi$  with animation parameter  $t$ .

(d) Calculate the spatial  $L^1$  norm

$$I(t) = \int_{x \in \mathbb{R}} \Phi(x, t)$$

of the fundamental solution. *Hint(s):* Note that  $I(t) = 2J(t)$  where

$$J(t) = \int_0^\infty \Phi(x, t) dx.$$

Calculate  $J(t)^2$ . Use  $y$  as a spatial variable of integration in one of the factors  $J(t)$ . Write what you get as an iterated integral and then as an integral of a function of two variables over the first quadrant. Use polar coordinates.

(e) How could you modify  $\Phi$  so that it satisfies  $u_t = ku_{xx}$  for non-unitary conductivity? *Hint(s):* Remember Problem 1 from Assignment 3 = Exam 1 about scaling in space and time. Also read section 10.4 of Haberman and see what Haberman defines as the fundamental solution.

(f) Given  $u_0 : \mathbb{R} \rightarrow \mathbb{R}$  with  $u \in C^0(\mathbb{R})$ , the function

$$u(x, t) = \int_{\xi \in \mathbb{R}} \Phi(x - \xi, t) u_0(\xi)$$

is called the **spatial convolution** of the fundamental solution with  $u_0$ . Show that this spatial convolution satisfies the initial value problem

$$\begin{cases} u_t = u_{xx} & \text{on } \mathbb{R} \times (0, \infty) \\ u(x, 0) = u_0(x), & x \in \mathbb{R} \end{cases}$$

for the heat equation on the whole real line.

(g) (Bonus) How can you modify the one-dimensional fundamental solution of the heat equation to obtain the fundamental solution of the heat equation  $\Phi : \mathbb{R}^n \times (0, \infty) \rightarrow \mathbb{R}$  on (all of)  $\mathbb{R}^n$ ?

**Problem 3** (wave equation: transverse oscillation model; Haberman section 4.2)

In the transverse oscillation model for the wave equation (think of the wave equation as a model for the vibration of a guitar string) one may assume a tension  $T$  and a lineal density  $\rho$  for the string. Thus, the mass of a small portion of the string modeled by the graph of a function  $u : [0, L] \rightarrow \mathbb{R}$  over the interval  $[a, b] \subset (0, L)$  is approximately  $\rho(b - a)$ . Thus, Newton's second law of motion ( $F = Ma$ ) for the **transverse displacement**  $u$  gives (approximately)

$$Ma_{\text{transverse}} = \rho(b - a) u_{tt}(x^*, t) \approx F_{\text{transverse}}(b, t) + F_{\text{transverse}}(a, t) \quad (1)$$

where  $x^*$  is some point with  $x^* \in (a, b)$  and  $F_{\text{transverse}}$  is the component of the tension force orthogonal to the (equilibrium of the) string, i.e., the vertical component if the equilibrium of your string is horizontal.

- (a) Draw a picture of the graph of  $u$ , representing the position of the string above the interval  $(a, b)$  at time  $t$ . Draw your picture so that  $u_x(a, t)$  and  $u_x(b, t)$  are different, i.e., so that the string is curving over the interval  $[a, b]$ .
- (b) If  $T$  represents the tension in the string to the right and is always tangent to the string, use similar triangles to find the components  $F_{\text{transverse}}(a, t)$  and  $F_{\text{transverse}}(b, t)$  of  $T$  at  $x = a$  and  $x = b$  respectively.
- (c) Substitute your values from part (b) into (1), divide by  $b - a$  and take the limit as  $b, a \rightarrow x$  to obtain the wave equation for  $u$ .

**Problem 4** (variable tension, transverse oscillations of a hanging chain; O'Neil Advanced Engineering Mathematics section 6.1)

- (a) If your derivation in Problem 3 assumed the tension  $T$  was constant, go back and derive the wave equation for transverse oscillations in which  $T = T(x)$  depends on the position  $x$ .
- (b) Use your model to find the fundamental modes of oscillation (separated variables solutions) for a hanging chain.
- (c) Animate some of your fundamental modes.

**Problem 5** (transverse oscillations of a rectangular drum, Haberman section 7.3)

- (a) Find the separated variables solutions of the boundary value problem

$$\begin{cases} u_{tt} = \Delta u, & \text{on } R \times [0, \infty) \\ u|_{\partial R} \equiv 0, & t > 0 \end{cases}$$

for the wave equation on the rectangle  $R = [0, L] \times [0, M]$ .

- (b) Choose specific positive values for  $L$  and  $M$ , and animate some of the fundamental modes you found in part (a).